

Phys 4062/5062 – Lecture Seven – Atom Field Interactions

Outline

Atom Field Interaction – Dipole Approximation

Time Dependent Evolution of Probability Amplitudes for Weak Fields

Atomic Response for Pulsed Excitation

Atom Field Interaction

Semi-Classical Treatment

- Electric field is classical
- Atom is a quantum mechanical entity

$$\text{Solve: } i\hbar \frac{\partial \psi}{\partial t} = \hat{H}\psi \quad (201)$$

$$\hat{H} = \hat{H}_0 + \hat{H}_1(t)$$

- H_0 is the unperturbed Hamiltonian
 - eigenvalues – energy levels
 - eigenfunctions – wavefunctions
$$\psi_n(r,t) = \psi_n(r)\exp[-iE_n t/\hbar] \quad (202)$$
 - ψ_n retains character of wavefunctions of H_0 such as parity and satisfies
$$\boxed{H_0\psi_n(r) = E_n\psi_n(r)}$$
- H_1 is the interaction of Electric Field that perturbs eigenfunctions of H_0

External Perturbation

$$\mathbf{E} = \mathbf{E}_0 \cos(\omega t) = |\mathbf{E}_0| \cos(\omega t) \hat{x} \quad (\text{Linearly polarized field})$$

$$H_1 = q\mathbf{r} \cdot \mathbf{E}_0 \cos(\omega t) = \mathbf{p} \cdot \mathbf{E}$$

$$\omega_0 = (E_2 - E_1)/\hbar \quad (\text{Resonant Frequency})$$

Introduce Rabi Frequency

$$\Omega = \frac{q\langle X_{12} | E_0 \rangle}{\hbar} \text{ where}$$

$$\langle X_{12} \rangle = \langle 1 | \hat{x} | 2 \rangle$$

$$\psi(r,t) = c_1 \psi_1(r) e^{-\frac{iE_1 t}{\hbar}} + c_2 \psi_2(r) e^{-\frac{iE_2 t}{\hbar}} \quad (\text{Linear Superposition of Unperturbed States})$$

$$\psi(r,t) = c_1 |1\rangle e^{-\frac{i\omega_1 t}{\hbar}} + c_2 |2\rangle e^{-\frac{i\omega_2 t}{\hbar}} \quad (204)$$

Here the probability amplitudes satisfies the normalization condition

$$|c_1(t)|^2 + |c_2(t)|^2 = 1 \quad (205)$$

Using (204) and (201)

$$i\dot{c}_1 = \Omega \cos(\omega t) e^{-i\omega_0 t} c_2 \quad (207)$$

$$i\dot{c}_2 = \Omega^* \cos(\omega t) e^{i\omega_0 t} c_1$$

Assume Initial Conditions

$$\begin{aligned} c_1(0) &= 1 \\ c_2(0) &= 0 \\ t \ll T &\Rightarrow |c_2(t)|^2 \ll 1 \end{aligned}$$

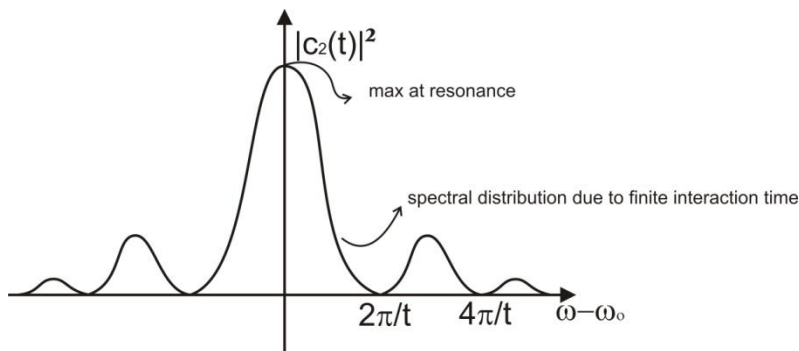
- there are weak field conditions
- T is the interaction time
- Integrate (207) to get

$$c_2(t) = \left(\frac{\Omega}{2}\right) \left[\frac{1 - \exp[i(\omega_0 + \omega)t]}{\omega_0 + \omega} + \frac{1 - \exp[i(\omega_0 - \omega)t]}{\omega_0 - \omega} \right] \quad (210)$$

- for optical frequencies $\omega \sim \omega_0$
- $|\Delta| = |\omega_0 - \omega| \ll \omega_0$
- $\omega_0 + \omega \sim 2\omega_0$
- this is called the rotating wave approximation which results in elimination of 1st term in (210)

Probability of Atom in $|2\rangle$ at time t is given by

$$|c_2(t)|^2 = |\Omega|^2 \left[\frac{\sin\left(\frac{(\omega_0 - \omega)t}{2}\right)}{\omega_0 - \omega} \right]^2 \quad (211)$$



- analogous to single diffraction pattern
- spread decreases when interaction time t increases just like diffraction angle decreases when slit width increases
- contrast this behaviour to steady state response
- relate atomic response to spectral components (Fourier components) associated with a square pulse